

Diffusion in microchannels II. Experimental Flow Characterization Electrofluidics.

Advection-Diffusion Equation



Advection-Diffusion equation



• D=0: concentration remains the same along a trajectory



• $D\neq 0$: concentration spreads with time

Taylor dispersion in pressure-driven flow

 material transported in a pressure-driven laminar flow will be spread due to diffusion and flow profile



Observation: sample band width changes is proportional to the square root of time, but with "effective" diffusion coefficient:

$$\frac{\partial c}{\partial t} = D(1 + \frac{Pe^2}{48})\frac{\partial^2 c}{\partial x^2} \quad \text{where Peclet number} \quad Pe = \frac{U_0 a}{D}$$

Modeling of time-dependent flow

- Model: spreading of a step distribution
- Parameters: D=1e-5; u=0.01 m/s



Advection-Diffusion equation

- Boundary conditions:
 - Dirichlet condition: c=0 at the wall



- homogeneous Neumann condition: $\partial c/\partial n=0$ at the wall absence of mass flux to the wall.
- Neumann condition: $J_n = -D\frac{\partial c}{\partial n} = \frac{d\Gamma}{dt} \quad \text{reaction kinetics}$

Coupling with Hydrodynamics

• The advection-diffusion equation requires knowledge of the velocity field; the complete set of equations:

$$\frac{\partial \rho}{\partial t} + \nabla \left(\rho \vec{U} \right) = 0$$
$$\rho \frac{\partial \vec{U}}{\partial t} + \rho \vec{U} \cdot \nabla \vec{U} = -\nabla P + \eta \Delta \vec{U} + F$$
$$\frac{\partial c}{\partial t} + \vec{U} \cdot \nabla c = \nabla \cdot \left(D \nabla c \right) + S$$

the system contains: 5 unknowns (u, v, w, P, c), 2 fluid properties and 2 external actions (F and S).

Peclet number

• diffusion advection equation without source or sink:

$$\frac{\partial c}{\partial t} + u \frac{\partial c}{\partial x} + v \frac{\partial c}{\partial y} + w \frac{\partial c}{\partial z} = D \left[\frac{\partial^2 c}{\partial x^2} + \frac{\partial^2 c}{\partial y^2} + \frac{\partial^2 c}{\partial z^2} \right]$$

the equation has 4 parameters: U_∞, L,c₀ and D.
 From Pi theorem, there is one nondimensional number:

$$Pe = \frac{U_{\infty}L}{D}$$

Distance of capture in a capillary

• by comparing axial convection to radial diffusion



$$\tau \approx \frac{R^2}{4D} \qquad L \approx 2V\tau \approx \frac{VR^2}{2D} \qquad L \approx \frac{Q}{2\pi D}$$
$$\frac{L}{R} = \frac{1}{2}Pe$$

Determination of diffusion coefficient

Microfluidic systems are convenient to determine diffusion coefficients in liquid phase



Mixing of fluids

- High laminarity of microflow delays mixing of the components
- to benefit advantages of microfluidics (e.g. reduction of reaction rates) mixing should be accelerated
- let's estimate a distance L at which a relative concentration on the wall is 90% of $c_0/2$.

$$c(x, y) = \frac{1}{2}c_0 \left(1 - erf \frac{y\sqrt{U}}{\sqrt{4Dx}}\right) \qquad \qquad \frac{R\sqrt{U}}{\sqrt{4DL}} = erfinv(0.1)$$
$$\frac{L}{R} \approx 32\frac{UR}{D} = 32Pe$$

Concentration boundary layer



 $L_{mix} \approx 32R \cdot Pe$ $h \approx 0.04 \cdot 2R \cdot \text{Re}_{D}$

- hydrodynamic entrance: the action of viscosity homogenizes the flow to reach a fully developed flow. Very short in microfluidics.
- mixing: the action of molecular diffusion homogenizes concentration

Improving the mixing

 as mixing length varies as square of R, dividing flow into thin brunches (achieved by lamination and hydrodynamic focusing) improves mixing



Mixing in two phase flow

Direction of plug motion





Mixing in digital microfluidics

 mixing in digital microfluidics shows very special pattern



Separation/purification of bioparticles



• Field Flow Fractionation (FFF): particles in a liquid flow separate according to their physical properties.

Chromatography



 Chromatography devices include two phases: mobile (transporting the sample) and stationary. Stationary phase is designed to retain longer the phase it associates with (e.g small fractions in size exclusion chromatography)

Proteomic reactor



Proteomic reactor



Microfluidics: Experimental Flow Characterization

Why knowing flow field is important

 Microfluidics devices used in variety if applications, e.g. Inkjet nozzles, medical microfluidic devices for diagnostics, PCR based DNA amplification/detection, channels for DNA fractinations, microtrusters on satellites and microaircrafts

• complicated geometries, flow difficult to simulate numerically

- Pointwise methods
 - Laser Doppler velocimetry (LDV). Lateral x longitudinal resolution down to 5um x 10 um
 - Optical Doppler tomography: suitable for highly scattered medium (e.g. Blood flow under the skin), resolution down to 5um x 15 um



- Full-Field techniques: generate two component velocities distributed in 2D plane. Essential for microfluidics
 - Scalar Image Velocimetry (SIV): uses molecular tracers to visualize the flow. Can go through smallest channels, insensitive to electric forces, but subject to intensive diffusion. Variation: dye uncaging





 Molecular tagging velocimetry: fluorescent/phosphorescent dye excited in a patteren, evolution of the pattern followed



from: Thijs Elenbaas, PhD Thesis, Univ. Eindhoven '06

- Particle Streak Velocimetry
- Particle Image Velocimetry (PIV)

Advantages:

- high spacial resolution
- large range of velocities covered (up to 8m/s)
- simplicity



Technique	Author	Flow Tracer	Spatial	Observation
			Resolution (µm)	
LDA	Tieu et al. (1995)		5 × 5 × 10	4–8 fringes limits velocity resolution
Optical Doppler tomography (ODT)	Chen et al. (1997)	1.7-µm polystyrene beads	5 × 15	Can image through highly scattering media
Optical flow using video microscopy	Hitt et al. (1996)	5-µm blood cells	$20 \times 20 \times 20$	In vivo study of blood flow
Optical flow using X-ray imaging	Lanzillotto et al. (1996)	1 – 20-µm emulsion droplets	~ 20 - 40	Can image without optical access
Uncaged fluorescent dycs	Paul et al. (1997)	Molecular Dye	$100 \times 20 \times 20$	Resolution limited by molecular diffusion
Particle streak velocimetry	Brody et al. (1996)	0.9-μm polystyrene beads	~ 10	Particle streak velocimetry
PIV	Urushihara et al. (1993)	1-µm oil droplets	$280\times280\times200$	Turbulent flows
Super-resolution PIV	Keane et al. (1995)	1-µm oil droplets	$50 \times 50 \times 200$	Particle tracking velocimetry
Micro-PIV	Santiago et al. (1998)	300-nm polystyrene particles	6.9 × 6.9 × 1.5	Hele-Shaw Flow
Micro-PIV	Santiago et al. (1998)	300-nm polystyrene particles	6.9 × 6.9 × 1.5	Silicon microchannel flow
Micro-PIV	Meinhart et al. (1999)	200-nm polystyrene particles	5.0 × 1.3 × 2.8	Microchannel flow

Comparison of High-Resolution Velocimetry Techniques [22].

Micro Particle Image Velocimetry

- Particle Image Velocimetry (PIV):
 - flow is seeded by small particles,
 - consecutive photographs of particles distribution are made
 - Images are sectioned into *interrogation regions*
 - Motion of particles within interrogation region determined by image cross-correlation

Micro Particle Image Velocimetry







$$\Phi(m,n) = \sum_{j=1}^{q} \sum_{i=1}^{p} f(i,j) \cdot g(i+m,j+n)$$



Physics consideration of micro-PIV

- Particles small considered to λ. If d<<λ, amount of light scattered varies as d⁶.
- Brownian motion

Flow particle dynamics

• Response time of the paricle

$$\tau_p = \frac{d_p^2 \rho_p}{18\eta}$$

- E.g. for 300nm polystyrene in water, approx 10⁻⁹ s
- Example: calculate responce time for polystyrene particles (ρ=1.05 g/cm³), diameters 1um, 300nm, 100nm in water (μ=10⁻³ kg/m s)

Velocity errors

- Particle is subject to Brownian motion
- High frequency velocity components of Brownian motion are not resolved, velocity meant as average displacement (after many fluctuations) per given time

$$\langle x \rangle^2 \propto D\Delta t$$

For spherical particle: $D = \frac{KT}{3\pi\mu d_p}$

Probablilty of particle displacement is described by 3D Gaussian

$$p(x, y, z) = \frac{\exp[-(x^2 + y^2 + z^2)/4D\Delta t]}{(4\pi D\Delta t)^{3/2}}$$

PIV measurements are done in 2D within some measurement depth:

$$p(x, y) = \frac{\exp[-(x^2 + y^2)/4D\Delta t]}{4\pi D\Delta t}$$

That's an approximation as the diffusion is still 3D and a particle can leave the measurement volume in z direction. Valid if $w\Delta t < \delta z_m$.

Brownian motion causes particle trajectories to fluctuate about the ideal pathlines.

Imagine ideal streamline followed for a period Δt :

$$\Delta x = u \Delta t$$
$$\Delta y = v \Delta t$$

Relative errors due to Brownian motion:

$$\varepsilon_{x} = \frac{\sigma_{x}}{\Delta x} = \frac{1}{u} \sqrt{\frac{2D}{\Delta t}}$$
$$\varepsilon_{y} = \frac{\sigma_{y}}{\Delta y} = \frac{1}{v} \sqrt{\frac{2D}{\Delta t}}$$

More significant for short measurement time slow velocities.



Positional uncertainty due to Brownian motion: particle displacement during the exposure time

Example: Compute dffusion coefficient D for 300nm polystyrene particle in water at room temperature. Given a flow speed of 10mm/s and 10 particles in an interroagation region, what Δt is required to achieve 10% relative error? 1% relative error? What particle smearing would you expect at 10ms exposure time?

• Density of polystyrene ρ =1.05 g/cm³,

•Viscosity of water $\mu = 10^{-3}$ kg/m s

3D diffraction pattern



2D Airy pattern



Depth of field

• Defined as $\frac{1}{4}$ of the distance between two minima, u=+- π



Depth of correlation

• The distance away from focal plane where intensity of the particle considered to be sufficiently small (usually 0.1 of in focal intensity)

$$\delta z_m = \frac{3n\lambda_0}{NA^2} + \frac{2.16d_p}{\tan(\theta)} + d_p$$

Example: what depth of correlation would you expect for x40 lens, 0.7 NA, illumination at 600nm for 1um, 300nm and 100nm particles.

Background noise

 Mainly due to particles out of focus.
 Can be lowered by reducing concentration of particles and depth of filed

Depth (µm)	Particle Concentration (by Volume)				
	0.01%	0.02%	0.04%	0.08%	
25	2.2	2.1	2.0	1.9	
50	1.9	1.7	1.4	1.2	
125	1.5	1.4	1.2	1.1	
170	1.3	1.2	1.1	1.0	

The Effect of Background Noise on Image Quality (i.e., Signal-to-Noise Ratio) [23]

Processing methods for micro-PIV

• Overlapping of low image density (LID) PIV: single image pair might exhibit noise, it will be averaged out if more images are considered



Correlation based overlapping

$$\Phi(m,n) = \sum_{j=1}^{q} \sum_{i=1}^{p} f(i,j) \cdot g(i+m,j+n)$$

In steady laminar flow, the peak position in correlation should be the same. Therefore: $1 \sum_{n=1}^{N} 1$

$$\Phi_{ens}(m,n) = \frac{1}{N} \sum_{k=1}^{N} \Phi_k(m,n)$$



Averaged on 101 image pair

Single image pair

- Background removal
 - Averaging background over many images
 - or
 - Using minimum value for every pixel in sequence



Micro-PIV examples

Measurements in microchannel

Equipment:

- 30um x 300um x 25mm glass channel
- •Epifluorescent microscope with x60 lens, NA=1.4. Focal plane 7.5um above the bottom of the channel
- light source: pulsed NdYAG laser
- Syringe pump. Flow rate 200ul/h
- 200nm fluorescent polystyrene particles



Interrogation window size: 120 x 8 pixels near the wall (13.6um x 0.9 um), 120 x 40 pixels everywhere else (13.6 x 40 um)



Results:

Measured and analytical solution match perfectly
possible to measure velocity profile down to within 450 nm of the wall
position of the wall can be determined within 8 nm



Flow in micronozzle



Flow in a biochip





Extension of micro-PIV

• Microfluidic "Nanoscope"

by careful inspection of velocity profiles wall position can be obtained with high precision (<100nm)



Useful to check channel geometry and hydrogel expansion

Micro-PIV thermometry

• Displacement (cross-correlation shift) and temperature (peak width) can be obtained independently



Particle tracking velocimetry or "super-resolution" PIV



PIV resolution can be increased if after first PIV step the average velocity is used to track individual particles. In this case more than 5200 velocity measurements were performed using 10 image pairs.

Electric field in a continuous media: Polarization

 Non-constant electric field will cause polarization of the media:



$$F_{i} = \int_{\Omega} d\vec{r} \rho(\vec{r_{0}} + \vec{r}) \left[E_{i}(\vec{r_{0}}) + r_{j}\partial_{j}E_{i}(\vec{r_{0}}) \right] = QE_{i}(\vec{r_{0}}) + p_{j}\partial_{j}E_{i}(\vec{r_{0}})$$
$$Q = \int_{\Omega} d\vec{r} \rho(\vec{r_{0}} + \vec{r})$$
$$p = \int_{\Omega} d\vec{r} \rho(\vec{r_{0}} + \vec{r})\vec{r}$$

Electrokinetic effects

- **Electrophoresis:** movement of charged surface (e.g. particle or molecule) relative to a stationary liquid.
- Electroosmosis: movement of liquid relative to a stationary charged surface (e.g. capillary tube)
- <u>Sedimentation potential</u>: electric potential created by moving charged particles
- <u>Streaming potential</u>: electric potential created by liquid moving relative to charged surface
- **Dielectrophoresis**: movement of uncharged with polarizability different from the liquid.

Debye layer in an electrolyte



• The potential at the Stern layer next to the surface is called zeta-potential ζ .

Electroosmotic flow

• Experimentally, if electric field is applied across a tube filled with electrolyte a uniform flow is induced:



Origin of electroosmosis

• Structure of the double layer



Origin of electroosmosis

• Flow velocity profile between parallel plates



Measuring of zeta potential

- tracking electroosmotic velocities
- streaming potential:
 - if a pressure driven flow is applied to an electrolyte, a potential will build up

$$\Delta V = \frac{\varepsilon \zeta}{\sigma \mu} \Delta p$$

Helmholtz-Smoluchowski equation

Electrophoretic effect

 Electrophoresis: motion of electrically charged molecule of particles in response to electric field Electrophoretic velocity



Dielectrophoresis

 Dielectrophoresis: movement of a charge neutral particle in a fluid induced by an inhomogeneous electric field (no DC field is necessary)



Particle is less polarizable than the fluid





Dielectrophoretic trapping of nanoparticles and molecules

 strong gradient of electric field existing near a sharp electrode will attract small polarizable particles/molecu les



S. Tuukkanen, J.J. Toppari, A. Kuzyk, L. Hirviniemi, V.P. Hytönen, T. Ihalainen, and P. Törmä, Nano Lett. **6**, 1339 (2006)).

Dielectrophoretic trapping of nanoparticles and molecules

 Metal nanoparticle trapped between the electrodes



A. Bezryadin, C. Dekker, and G. Schmid, Appl. Phys. Lett. 71, 1273—1275 (1997).

 Measurement of electrical characteristic of trapped polyC-polyG DNA

Danny Porath, Alexey Bezryadin, Simon de Vries & Cees Dekker, Nature 403, 635 (2000).





Electrokinetic valve

• Geometry



Electrokinetic valve

