

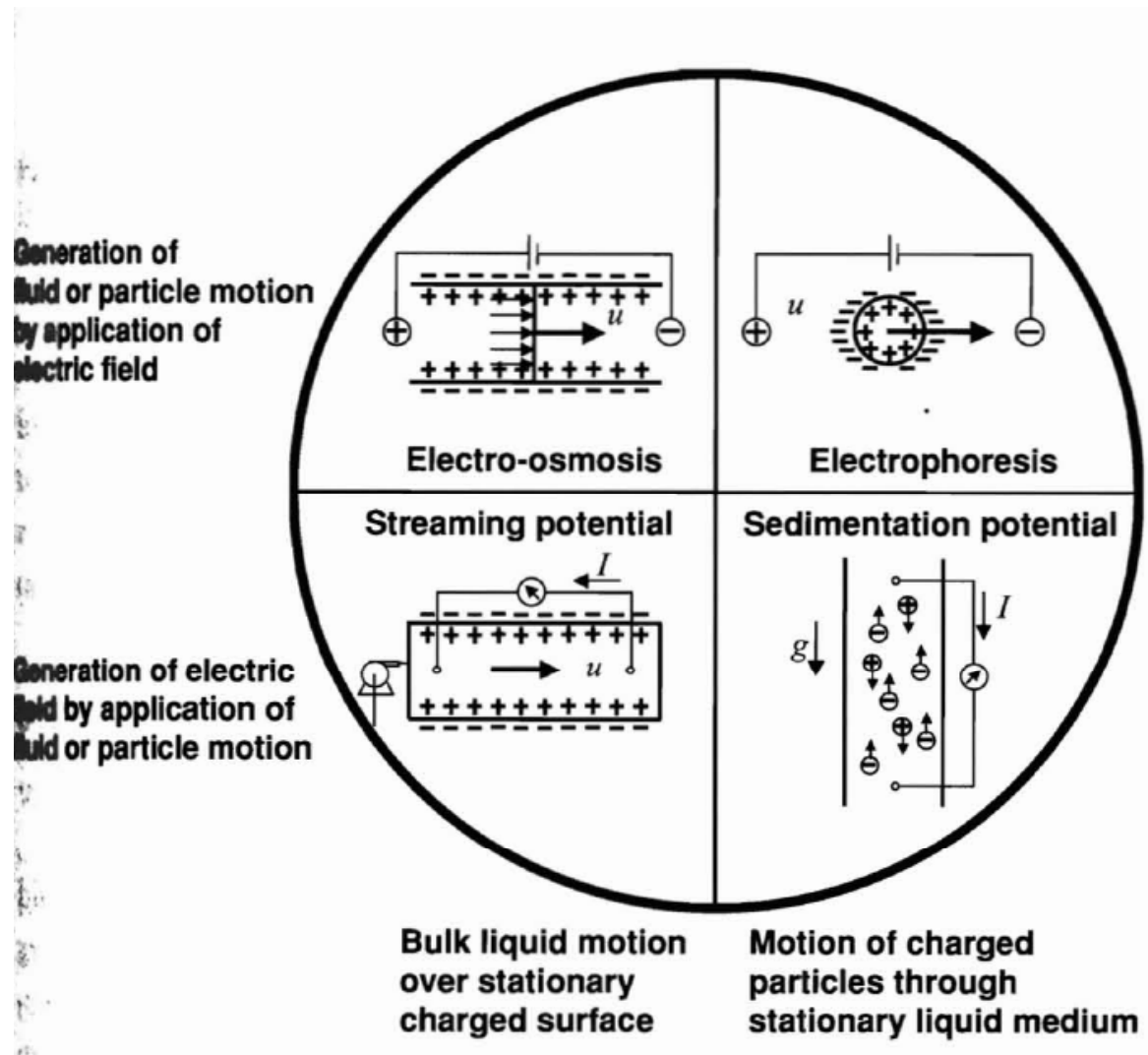
Lecture 8

Electrokinetic effects.
Electroosmosis and
electrophoresis.

Electrokinetic effects

- **Electrophoresis**: movement of charged surface (e.g. particle or molecule) relative to a stationary liquid.
- **Electroosmosis**: movement of liquid relative to a stationary charged surface (e.g. capillary tube)
- **Sedimentation potential**: electric potential created by moving charged particles
- **Streaming potential**: electric potential created by liquid moving relative to charged surface
- **Dielectrophoresis**: movement of uncharged particles with polarizability different from the liquid.

Electrokinetic effects



Electrophoresis

- Let's consider a charged spherical particle in a low conductivity liquid.

$$F_{drag} = 6\pi\eta a \vec{s} \quad \vec{F}_{el} = Ze \cdot \vec{E}$$

- the drift velocity is:

$$\vec{u}_{ep} = \frac{Ze \cdot \vec{E}}{6\pi\eta a} = \mu_{ion} \cdot \vec{E}$$

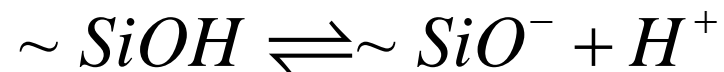
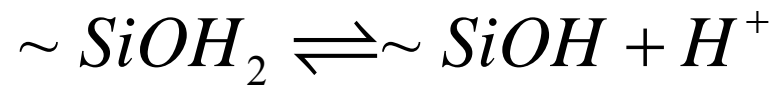
- the equation is valid also for ions (except H^+ and OH^-) assuming hydration radius of approx. 0.2nm (one layer of water molecules).

Origin of electroosmosis

- Most of the materials carry some charge on the surface (due to adsorption, presence of partially dissociated groups etc.)
- This charge will be screened by the ions in the solution
- If an external voltage is applied the ions will start to move in the electric field and cause the flow of the liquid

Charge on Si surface

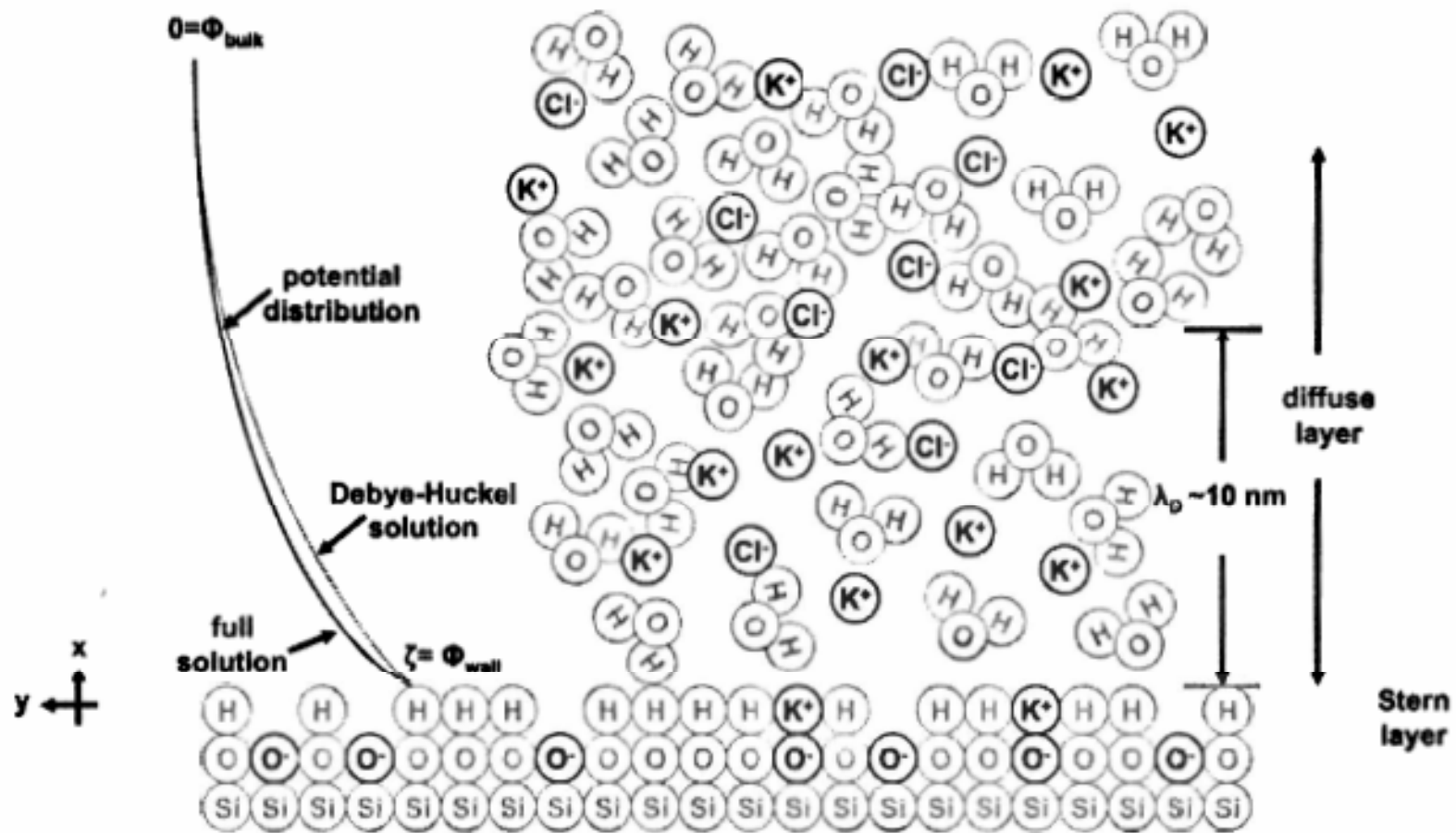
- In the case of oxides the dominating mechanism of surface charging is the dissociation of hydroxyl groups on the surface



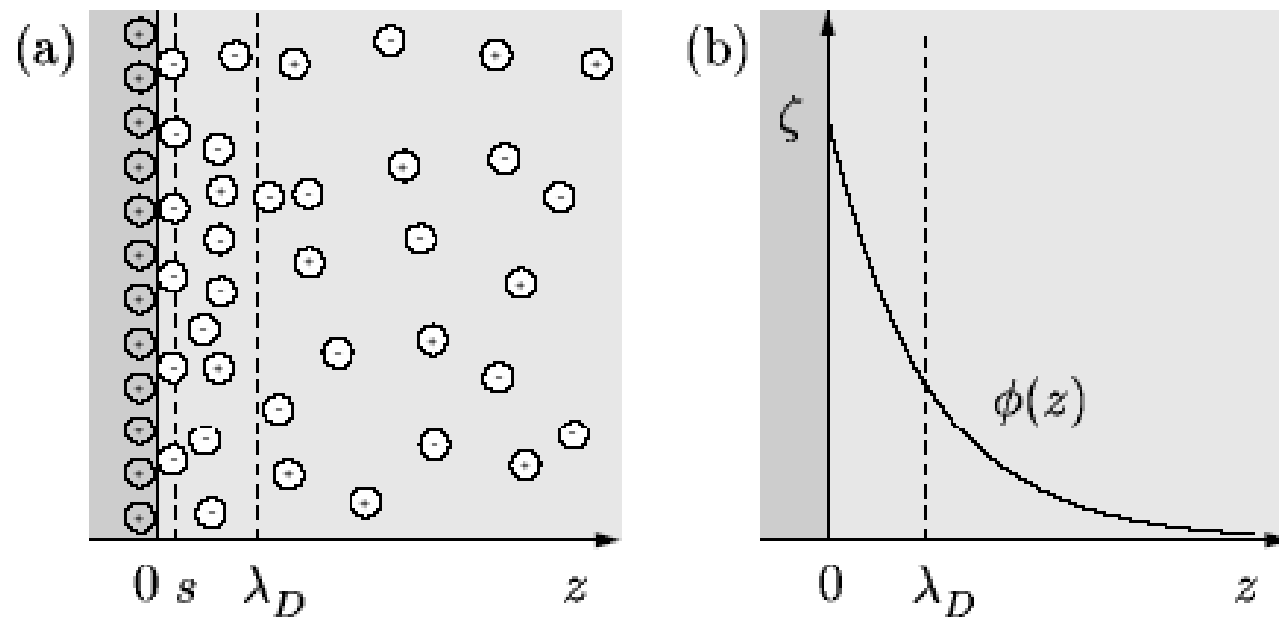
pzc ~ pH 1.8 - 3.4

Origin of electroosmosis

- Ionic structure of the solid (silicon) - liquid (electrolyte) interface



Debye layer in an electrolyte



- The potential at the Stern layer next to the surface is called zeta-potential ζ .

Debye layer

- We can find the charge and potential distribution by linking thermodynamics and electrodynamics:
 - The equilibrium condition: all ions have the same chemical potential:

$$\mu_{\pm}(r) = \mu_0(r) + kT \ln \left[\frac{c_{\pm}(r)}{c_0} \right] \pm Ze\varphi(r); \nabla \mu_{\pm}(r) = 0$$

$$kT \nabla \ln \left[\frac{c_{\pm}(r)}{c_0} \right] = \mp \nabla Ze\varphi(r)$$

the boundary conditions: $c_{\pm}(\infty) = c_0$; $\varphi(\infty) = 0$

$$c_{\pm}(r) = c_0 \exp \left[\mp \frac{Ze}{kT} \varphi(r) \right]$$

Debye layer

- Now we can link charge distribution and it's electrical potentials using Poisson equation

$$\left\{ \begin{array}{l} \nabla^2 \varphi(r) = -\frac{\rho_e(r)}{\varepsilon} \\ \rho_{el} = Ze[c_+(r) - c_-(r)] = -2Zec_0 \sinh\left[\frac{Ze}{kT}\varphi(r)\right] \end{array} \right.$$

- Poisson-Boltzmann equation

$$\nabla^2 \varphi(r) = 2\frac{Zec_0}{\varepsilon} \sinh\left[\frac{Ze}{kT}\varphi(r)\right]$$

Debye-Hückel approximation

- In the limit of small zeta-potential: $Ze\zeta \ll kT$

$$\nabla^2 \varphi(r) = 2 \frac{(Ze)^2 c_0}{kT \varepsilon} \varphi(r) = \frac{1}{\lambda_D^2} \varphi(r)$$

$$\lambda_D = \sqrt{\frac{\varepsilon kT}{2(Ze)^2 c_0}} \approx \sqrt{\frac{1mM}{Z^2 c_0}} 9.6nm$$

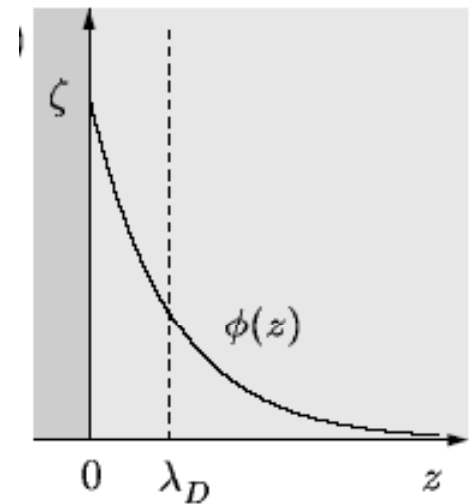
Debye-Hückel approximation: an infinite wall

- 1-D situation of a potential/charge distribution next to an infinite wall

$$\nabla^2 \varphi(z) = \frac{1}{\lambda_D^2} \varphi(z)$$

the solution should be in the form: $\varphi(z) = C_1 e^{-z/\lambda_D} + C_2 e^{z/\lambda_D}$

the boundary conditions: $\varphi(\infty) = 0$ $\varphi(0) = \zeta$



solution:

$$\varphi(z) = \zeta \exp\left[-\frac{z}{\lambda_D}\right]$$

$$\rho_{el}(z) = -\varepsilon \frac{\partial^2 \varphi}{\partial z^2} = -\frac{\varepsilon \zeta}{\lambda_D^2} \exp\left[-\frac{z}{\lambda_D}\right]$$

Debye-Hückel approximation: parallel plate channel

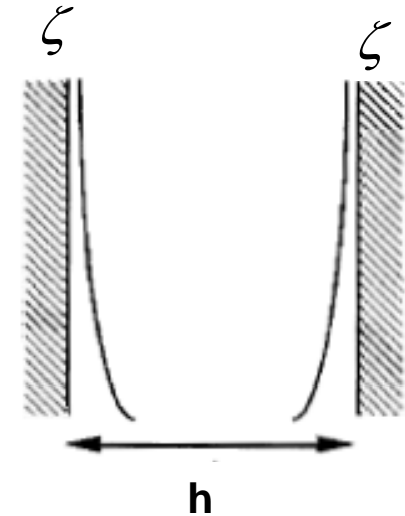
- 1-D situation of a potential/charge distribution in a parallel plate channel

$$\nabla^2 \varphi(z) = \frac{1}{\lambda_D^2} \varphi(z)$$

the solution should be in the form: $\varphi(z) = C_1 e^{-z/\lambda_D} + C_2 e^{z/\lambda_D}$

the boundary conditions:

$$\varphi(-h/2) = \zeta \quad \varphi(h/2) = \zeta \quad \left. \frac{\partial \varphi}{\partial z} \right|_{z=0} = 0$$



solution:

$$\varphi(z) = \zeta \frac{\cosh[z/\lambda_D]}{\cosh[h/2\lambda_D]}$$

$$\rho_{el}(z) = -\varepsilon \frac{\partial^2 \varphi}{\partial z^2} = -\frac{\varepsilon \zeta}{\lambda_D^2} \frac{\cosh[z/\lambda_D]}{\cosh[h/2\lambda_D]}$$

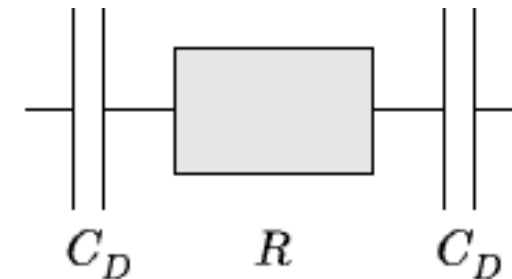
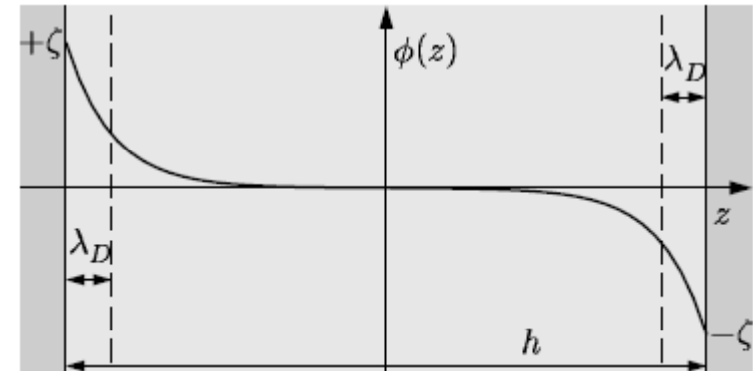
Capacitance and characteristic frequency of a double layer

- The Debye layer acts as an electrical capacitor
- The capacitance of the Debye layer:

$$q_{el} = \int_0^{\infty} \left[-\frac{\varepsilon \zeta}{\lambda_D^2} \exp\left(-\frac{z}{\lambda_D}\right) \right] dz = -\frac{\varepsilon}{\lambda_D} \zeta$$

$$C_D = \frac{\varepsilon}{\lambda_D}$$

$$\tau_{RC} = RC = \left(\frac{h}{\sigma_{el} A} \right) \left(\frac{1}{2} \frac{\varepsilon}{\lambda_D} A \right) = \frac{1}{2} \frac{h \varepsilon}{\sigma_{el} \lambda_D}$$

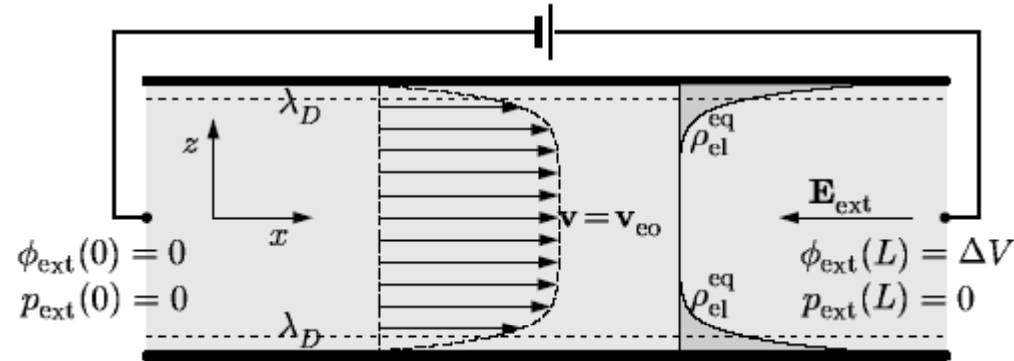


- Typically, for 1mM solution:

$$\lambda_D = 9.6 \text{ nm}; h = 100 \mu\text{m}, \sigma_{el} = 10^{-3} \text{ S/m}; \varepsilon = 78 \varepsilon_0 \quad \Rightarrow \quad \tau_{RC} = 3.6 \text{ ms}$$

Electroosmosis

- We have an extra body force in the Navier-Stokes equation due to electrical field:



$$\rho \left[\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \cdot \vec{v} \right] = -\nabla p_{ext} + \eta \nabla^2 \vec{v} - \rho_{el} \nabla \phi_{ext}$$

$$E_{ext} = -\nabla \phi$$

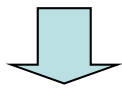
- Comments:
 - we assume that the external potential doesn't change the charge distribution
 - for continuous EO flow, some electrochemical processes on electrodes are required

Electroosmosis

$$\rho \left[\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \cdot \vec{v} \right] = -\nabla p_{ext} + \eta \nabla^2 \vec{v} - \rho_{el} \nabla \phi_{ext}$$

$$\vec{v} = (v_x(z), 0, 0)$$

$$\nabla p_{ext} = 0$$

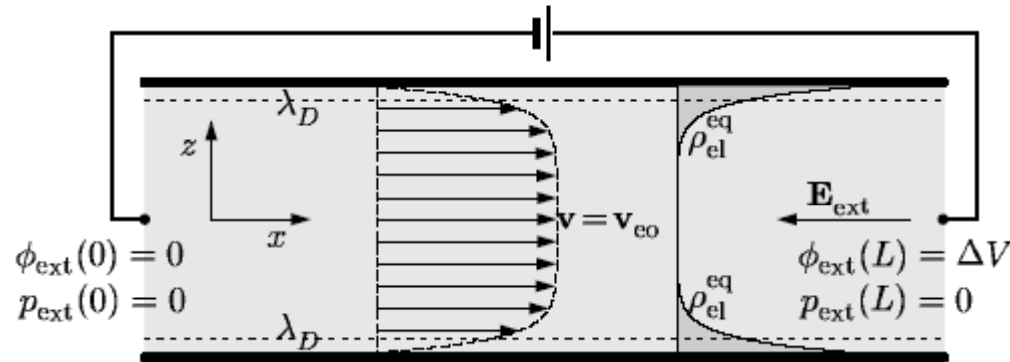


$$0 = \eta \frac{\partial^2 v_x}{\partial z^2} + \varepsilon \frac{\partial^2 \phi_{eq}}{\partial z^2} E \quad \text{or} \quad \frac{\partial^2}{\partial z^2} \left[v_x(z) + \frac{\varepsilon E}{\eta} \phi(z) \right] = 0$$

- the boundary conditions: $v_x(-h/2) = v_x(h/2) = 0$

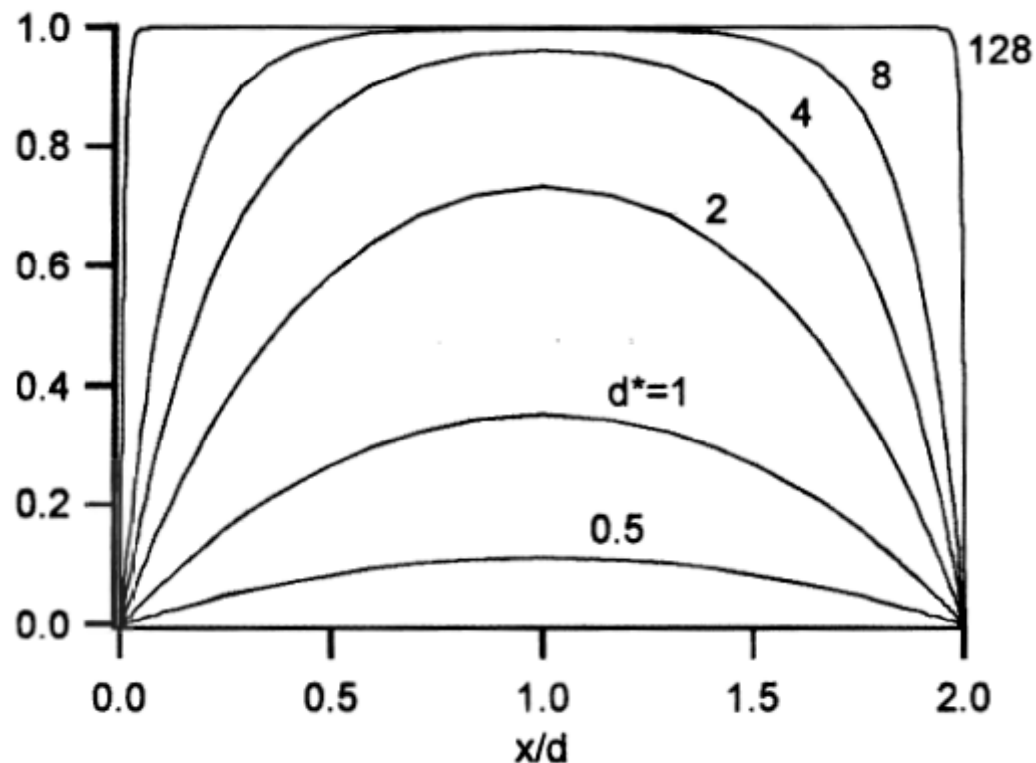
- the solution
$$v_x(z) = \left[\zeta - \phi_{eq}(z) \right] \frac{\varepsilon E}{\eta} = \frac{\varepsilon \zeta}{\eta} E \left[1 - \frac{\cosh[z/\lambda_D]}{\cosh[h/2\lambda_D]} \right]$$

v_{EO}



Electroosmosis

- Flow velocity profile between parallel plates

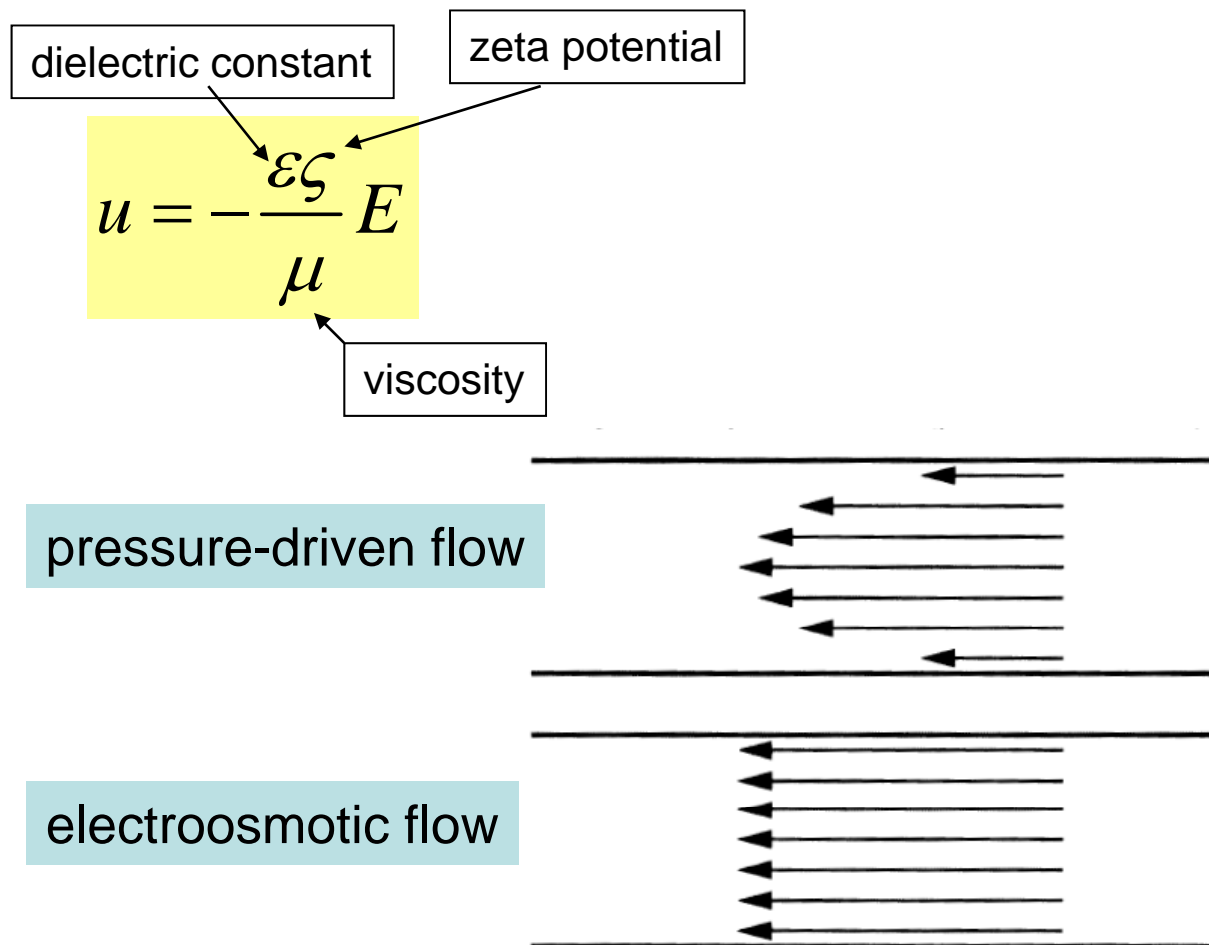


$$u^* = u (\varepsilon \zeta E / \mu)^{-1}$$

$$d^* = d / \lambda_D$$

Electroosmotic flow

- typical values are: $\zeta \approx 100\text{ mV}$; $\lambda_D \approx 10\text{ nm}$;
- As we see, if electric field is applied across a tube filled with electrolyte a fairly uniform plug flow is produced (vs. parabolically pressure driven)



Measuring of zeta potential

- tracking electroosmotic velocities
- streaming potential:
 - if a pressure driven flow is applied to an electrolyte, a potential will build up

$$\Delta V = \frac{\varepsilon \zeta}{\sigma \mu} \Delta p$$

Helmholtz-Smoluchowski equation

Electrophoretic effect

- Electrophoresis: motion of electrically charged molecule of particles in response to electric field
Electrophoretic velocity

zeta potential at the particle surface

$$u = \frac{\varepsilon \zeta}{\eta} E$$

$$\lambda_d \ll d \quad (\text{particle size})$$

$$u = \frac{2}{3} \frac{\varepsilon \zeta}{\eta} E$$

$$\lambda_d \gg d$$

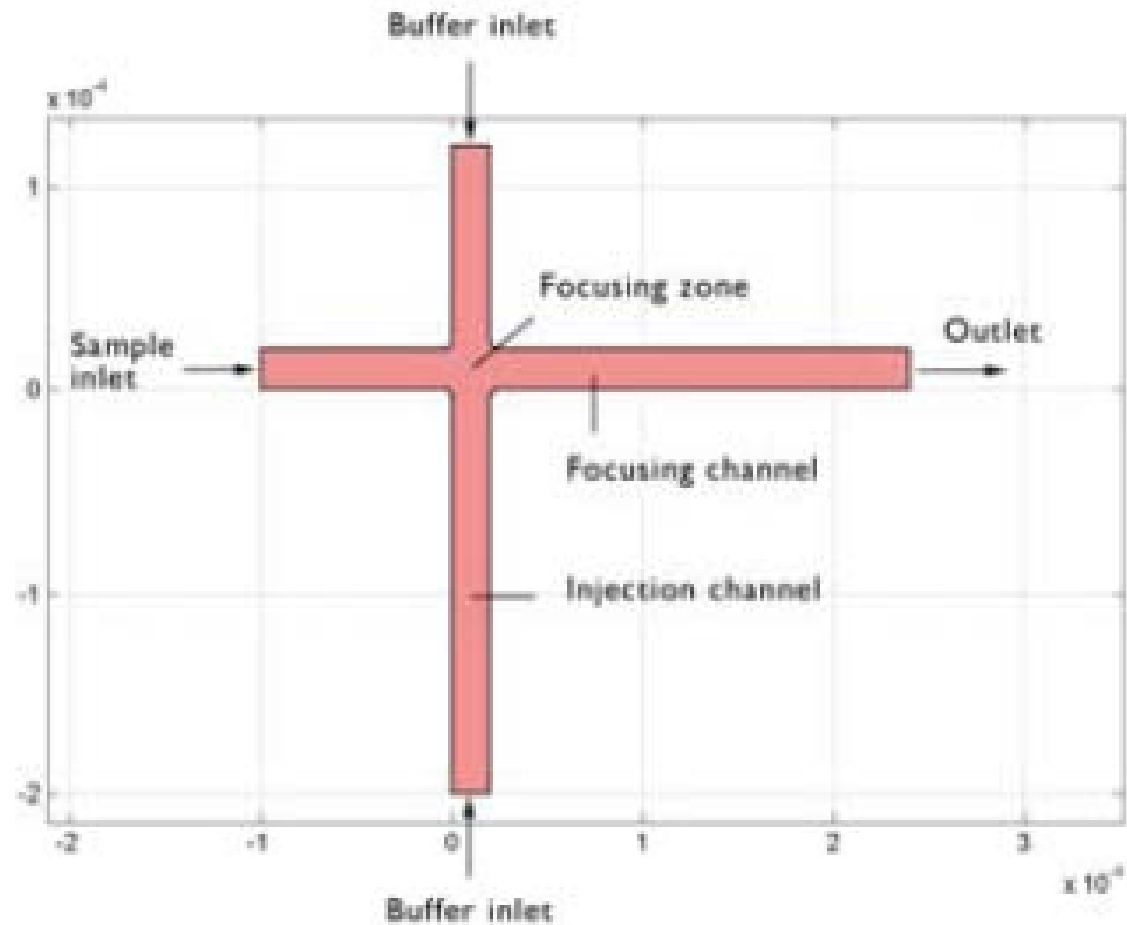
Generally: $\mu_{ep} = A \frac{\varepsilon \zeta}{\eta} E; \quad \mu_{net,j} = \mu_{eo} + \mu_{ep,j}$

Resolution of electrophoretic system: $R_{mn} = \frac{\Delta t_{mn}}{w}$

$\Delta t_{mn} \leftarrow \propto t$
 $w \leftarrow \propto \sqrt{t}$

Electrokinetic valve

- Geometry



Electrokinetic valve

- Technique

